

# Hearing what gravitational wave standard sirens have to say

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### Talk outline

- Motivation
- Current state of knowledge
- The main ideas in GW cosmology
  - GW "standard sirens"
  - Statistical arguments
  - Space-based approaches
- Some recent ideas regarding neutron stars
  - Using tidal signatures
  - Using the hyper-massive NS
- Summary

#### Motivation

### Gravitational waves

- Gravitational waves are propagating oscillations of the gravitational field.
- Travelling at the speed of light.
- Composed of 2 polarisations.

 Generated by time varying mass quadrupole (and higher) moment(s).



### Detection rates

#### LIGO-Virgo Collaboration, arXiv:1304.0670 (2013)

$\begin{array}{ c c c c c c c c c c c c c c c c c c c$		Estimated			Number	% BNS	Localized		
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$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	Epoch	Duration	LIGO	Virgo	Detections	$5{ m deg}^2$	$20\mathrm{deg}^2$		
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	2015	3 months	40 - 80	—	0.0004 - 3	—	_		
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	2016 - 17	6 months	80 - 120	20 - 60	0.006 - 20	2	5 - 12		
$\begin{array}{c c c c c c c c c c c c c c c c c c c $	2017 - 18	9 months	120-170	60 - 85	0.04 - 100	1 - 2	10 - 12		
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Table 5. Detection rates for compact binary coalescence sources.ruled out in 2015IFOSource <sup>a</sup> $\dot{N}_{low}$ yr <sup>-1</sup> $\dot{N}_{re}$ yr <sup>-1</sup> $\dot{N}_{high}$ yr <sup>-1</sup> $\dot{N}_{max}$ yr <sup>-1</sup> 2015NS-NS $2 \times 10^{-4}$ $0.02$ $0.2$ $0.6$ NS-BH $7 \times 10^{-5}$ $0.004$ $0.1$ InitialBH-BH $2 \times 10^{-4}$ $0.007$ $0.5$ IMRI into IMBH $2 \times 10^{-4}$ $0.007$ $0.01^{c}$ IMBH-IMBH $0.4$ $400$ $400$ $1000$ NS-NS $0.4$ $40$ $400$ $1000$ NS-BH $0.2$ $10$ $300$ $0.4$ AdvancedBH-BH $0.4$ $20$ $10000$ $1000$ IMRI into IMBH $0.4$ $20$ $1000$ $0.6$	2022 + (India)	(per year)	200	130	0.4 - 400	17	48		
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85000000		BH–BH	0.4	20	1000			different	
		IMRI into IMBH			10 <sup>b</sup>	300 <sup>c</sup>		assumptions	
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LIGO-Virgo Collaboration, CQG 27, 173001 (2010)

### Motivation

- GW detection *alone* will be \*pretty\* good, *but...*
- The detection and characterisation of a *population* of GW sources will allow
  - the study of the large-scale structure of the Universe.
  - us to infer the formation history of the massive black hole population.
  - precision mapping of the expansion history of the Universe.
  - the use of cosmic distance markers (standard sirens).
  - provide a "powerful" probe of the dark energy content of the universe.

#### The current state of knowledge

### Cosmic distance ladder



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- Nearby objects are used to calibrate more distant measurements.
- GW measurements would be **independent** of this ladder.

### Hubble diagram

• Redshift 
$$1+z = \sqrt{\frac{1+v/c}{1-v/c}}$$

• Luminosity distance

$$D_{\rm L} = c(1+z) \int_{0}^{z} \frac{dz'}{H(z')}$$

• Hubble parameter





### Standard Candles

- Type 1a supernovae progenitors are thought to be white dwarfs pushed over the Chandrasekhar limit.
- They act as standard candles" of \*equal\* luminosity (to ~15%).
- Calibration with Cepheids gives  $H_0 =$ 73.8 ± 2.4 km s<sup>-1</sup> Mpc<sup>-1</sup> [Riess *et al*, ApJ (2011)]



### Current knowledge

- The recently published Planck CMB results (combined with others) give the best constraints to date.
- Consistent with the standard ΛCDM model.
- Gives  $H_0 = 67.8 \pm 0.9$ km s<sup>-1</sup>Mpc<sup>-1</sup>
- These (EM) results are likely to improve before GWs are competitive.



#### GW standard sirens

### Standard Candles

- The inverse square law relates the received flux to the distance
- All you need to know is that there are events/objects of equal intrinsic luminosity e.g. Type 1a Supernova



Gives relative distances so still need calibration







### GW standard sirens

- The inverse square law relates the received GW amplitude to the distance
- GW compact-binary-coalescences of neutron stars or black holes (of given mass) will produce identical signals.



Gives absolute distances so no \*calibration\* required



### Schutz's idea

- Schutz in 86' proposed using compact-binary-coalescences as "standard sirens" [Schutz, Nature (1986)].
- Phase measurement gives redshifted chirp-mass  $\mathcal{M}_z = \mathcal{M}(1+z)$

- Amplitude gives ratio of redshifted chirp-mass<sup>5/3</sup> with luminosity distance D<sub>L</sub>.
- "Self-Calibrating" sources but **no redshift**.

### $\mathcal{M}, \mathcal{Z}$ degeneracy

The problem is that we only get D<sub>L</sub> and the redshifted mass

 $\mathcal{M}_z = \mathcal{M}(1+z)$ 

- We need EM measurements of redshift to break the degeneracy.
- Therefore we need host galaxy identification.



### Gamma-ray bursts

 GRBs represent an EM counterpart with redshift obtained from the host galaxy. [Dalal et al PRD (2006), Nissanke et al ApJ (2010), Zhao et al PRD (2011)]





### Schutz's method



### Galaxy catalogues

- Del Pozzo extended the Schutz idea to make use of galaxy catalogues to identify hosts [Del Pozzo PRD (2012)].
- The redshift can then be obtained.
- Any confusion on between host galaxies is averaged out with many sources.



### Statistical arguments

### Statistical properties

 Idea first proposed by Marković 93' and Finn & Chernoff 93' to use the distribution of measured SNRs. [Markovic PRD (1993), Finn & Chernoff ApJ (1993), Finn PRD (1996)]



### Statistical properties

- The idea was expanded upon by Taylor *et al* 2011[Taylor *et al* PRD (2011), Taylor *et al* PRD (2012)]
- Where the mass distribution and star formation rate are included in the model.



Taylor et al PRD (2012)

### Using tidal signatures

### Neutron star Equation of State

- The equation of state (EOS) specifies the pressure of neutron star matter at a given density.
- Expect a single equation of state to describe all (cold) neutron stars
- It's not crazy to think that the Neutron-star EOS will be well-understood in the era of third-generation GW detectors [Del Pozzo et al PRL (2013)].



### Tidal deformation

- Each neutron star's tidal field deforms the other star.
- The EOS sets how a neutron star of given mass responds.
- Tidal deformation modifies orbital energy and GW luminosity, contributing to the GW phase evolution:



$$\Phi(f)_{\text{tidal}} = \frac{3}{128\eta} (\pi M f)^{5/3} \left( -\frac{24}{M^5} \left( \frac{M+11m_1}{m_2} \lambda_1 + \frac{M+11m_2}{m_1} \lambda_2 \right) + \dots \right)$$

### Tidal deformation

- Each EOS will provide a different level of deformation.
- Hence a different level of GW phasing.
- Also a different dependence on NS mass
- Tidal deformability is a function of the Love number k<sub>2</sub>



$$\lambda(m) = (2/3)k_2R^5(m)$$

### EOS effect on waveforms

• The presence of matter modifies the late-inspiral, merger, and post-merger GW signals: the high-frequency part of the coalescence.



# Breaking M-z degeneracy

- Rest-frame waveform phase evolution with tidal contributions: (leading order, equal mass system)
- Tidal terms are formally 5 and 6PN order

$$\Phi(f) = \Phi_{\rm PN} - \frac{117}{8\eta} (\pi M (1+z)f)^{5/3} \frac{\lambda_1({\rm EOS};m)}{M^5}$$
  
These mass-dependent terms are not paired with a redshift factor

 Inference on the detected waveform allows us to estimate M and z.

### Redshift measurement



- For 3rd generation detectors the GW signal alone can be used to determine the redshift of the source.
- Even in the worst case, redshift uncertainties can be constrained to ~40% at z < 1 and in the best case ~8%.

### Cosmological implications



### Using the hyper-massive NS

### Post-Merger

- When 2 NSs merge after the inspiral they **briefly** form an unstable hyper-massive NS (if the EOS is not too soft).
- Such an object will survive for O(10s) milliseconds before collapsing into a black-hole.
- The GW waveform contains signatures of the EOS encoded in multiple frequency components.



Takami et al arXiv:1412:3240 (2014)

### Post-Merger waveforms



Takami et al arXiv:1412:3240 (2014)

### Post-Merger waveforms



### Dependence on mass

- The frequency features correlate positively with mass.
- This is expected since heavier NSs are smaller
   → characteristic frequencies are higher.

$$R \propto M^{-1/3}$$



This work uses 5 waveforms from the same polytropic EOS

### The idea

- First detect the inspiral and accurately measure the redshifted mass.
- Then analyse the postmerger signal and measure the redshifted frequency feature(s).
- Their different dependence on redshift allows us to break the degeneracy.



Messenger et al PRX (2014)

### Redshift measurements

- Nice but despite the huge SNR in the inspiral, the post-merger has low SNR.
- Hence currently only applicable to nearby sources (Adv detector distances).
- However, still independent of EM observations.



# Summary

- GW sources are (will be) very useful cosmological probes.
- They will provide measurements **independent** of the "cosmic distance ladder".
- We have a number of different methods with and without EM counterparts.
- Calibration may end up being a limiting systematic factor.
- We need to compare our potential sensitivities to future EM experiments.
- Focus **right now** is on first direct detection.

#### Thanks

#### Extra slides

### Super-Massive binary black-holes

- *D*<sub>L</sub>,*z* relation investigated for LISA by Holz & Hughes
   2005. [Holz & Hughes 2005 ApJ]
- Statistical approach taken by Petiteau *et al* 2011. [Petiteau *et al* 2011 ApJ]
- Good localisation makes host identification tractable.
- Weak gravitational lensing is a limiting factor in estimating luminosity distance.



### Expansion acceleration

- Directly measuring the expansion of the universe during a GW event [Seto et al 2001 PRL].
- Observe for long enough to see an object's changing redshift.



### NS tidal effects

- CM & Read discovered that tidal effects in NS binaries break the M,z degeneracy.
   [Messenger & Read 2012 PRL, Li et al 2013]
- The additional phase contribution is a function of the intrinsic mass!
- So you get the redshift without an EM observation.

